

Refer to <https://www.femm.info/wiki/AxisymmetricFormulation> for nomenclature. Copied here for convenience...

FEMM Axisymmetric Interpolation Functions and Flux Density

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Introduction

The purpose of this note is to explain the somewhat unusual formulation used in FEMM's axisymmetric formulation. This is basically the method for first-order triangles considered in Section IV of Henrotte *et al.* The basic difference between this and other first-order methods is that this method interpolates radially in terms of r^2 , rather than in terms of just r . The result is a formulation that is more physically consistent than many other competing methods, yielding an accurate solution close to $r = 0$ where other methods can break down.

Interpolation

Instead of the regular magnetic vector potential A , the nodal field values are in terms of nodal flux ϕ , where A and ϕ are related by:

$$\phi = 2\pi r A$$

For ease of notation, define:

$$s = r^2$$

Now, consider a "unit triangle" that lies between the points $(p, q) = (0, 0), (1, 0), (0, 1)$ as pictured in Figure 1.

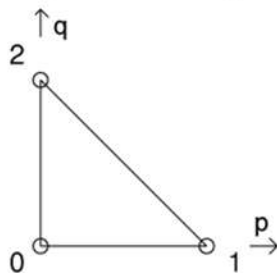


Figure 1: Unit triangle.

We can map the (p, q) coordinates back to the physical (s, z) coordinates using a standard linear interpolation scheme:

$$s = (s_1 - s_0)p + (s_2 - s_0)q + s_0$$

$$z = (z_1 - z_0)p + (z_2 - z_0)q + z_0$$

Nodal ϕ is interpolated in a similar way as:

$$\phi = (\phi_1 - \phi_0)p + (\phi_2 - \phi_0)q + \phi_0$$

This relationship between element and physical coordinates can be inverted to yield:

$$\begin{Bmatrix} p \\ q \end{Bmatrix} = \frac{1}{d} \begin{bmatrix} (z_2 - z_0) & -(s_2 - s_0) \\ -(z_1 - z_0) & (s_1 - s_0) \end{bmatrix} \begin{Bmatrix} s - s_0 \\ z - z_0 \end{Bmatrix}$$

where

$$d = (s_1 - s_0)(z_2 - z_0) - (z_1 - z_0)(s_2 - s_0)$$

We could write this more compactly as:

$$\begin{Bmatrix} p \\ q \end{Bmatrix} = \frac{1}{d} \begin{bmatrix} b_1 & c_1 \\ b_2 & c_2 \end{bmatrix} \begin{Bmatrix} s - s_0 \\ z - z_0 \end{Bmatrix}$$

where

$$d = c_2 b_1 - b_2 c_1$$

where the b 's and c 's are the corresponding entries in the previous matrices.

It is interesting to note that interpolating s rather than r produces an element with "warped" sides in the rz plane.

Flux Density

In terms of ϕ , the flux density can be written as:

$$B_r = -\frac{1}{2\pi r} \frac{\partial \phi}{\partial z} \quad B_z = \frac{1}{2\pi r} \frac{\partial \phi}{\partial r}$$

To be easily integrated with the fashion in which we are interpolating the element, the derivatives can be re-written as in terms of s rather than r as:

$$B_r = -\frac{1}{2\pi\sqrt{s}} \frac{\partial \phi}{\partial z} \quad B_z = \frac{1}{\pi} \frac{\partial \phi}{\partial s}$$

Writing out the derivatives of the interpolation function yields:

$$\begin{Bmatrix} \frac{\partial \phi}{\partial z} \\ \frac{\partial \phi}{\partial z} \end{Bmatrix} = \begin{Bmatrix} \frac{\partial \phi}{\partial z} \frac{\partial z}{\partial z} + \frac{\partial \phi}{\partial r} \frac{\partial z}{\partial z} \\ \frac{\partial \phi}{\partial z} \frac{\partial z}{\partial z} + \frac{\partial \phi}{\partial r} \frac{\partial z}{\partial z} \end{Bmatrix} = \frac{1}{d} \begin{bmatrix} b_1 & b_2 \\ c_1 & c_2 \end{bmatrix} \begin{Bmatrix} \phi_1 - \phi_0 \\ \phi_2 - \phi_0 \end{Bmatrix}$$

Substituting into the formulas for the derivatives yields:

$$B_r = -\frac{1}{2\pi d \sqrt{s}} (c_1(\phi_1 - \phi_0) + c_2(\phi_2 - \phi_0))$$

$$B_{rz} = -\frac{1}{\pi d} (b_1(\phi_1 - \phi_0) + b_2(\phi_2 - \phi_0))$$

Interpretation of the form of B

Now, we can observe that B_z is constant over the element, which is a reasonable behavior for a first-order element. The B_r flux becomes smaller as r increases, but the $1/r$ scaling factor is consistent with a constant amount of radially directed flux being spread over a larger area as the radius increases. This form is similar to the standard 3-node triangle formulation for 2-D planar problems, in which flux is constant over each element.

References

F. Henrotte *et al.*, "A new method for axisymmetric linear and nonlinear problems," *IEEE Transactions on Magnetics*, 9(2):1352-1355, March 1993. doi: 10.1109/CEFC.1992.720582*

In[1]:= **Br = -1 / (2 Pi r d) (c1 (phi1 - phi0) + c2 (phi2 - phi0))**

Out[1]=
$$-\frac{c_1 (-\phi_0 + \phi_1) + c_2 (-\phi_0 + \phi_2)}{2 d \pi r}$$

In[2]:= **Bz = 1 / (Pi d) (b1 (phi1 - phi0) + b2 (phi2 - phi0))**

Out[2]=
$$\frac{b_1 (-\phi_0 + \phi_1) + b_2 (-\phi_0 + \phi_2)}{d \pi}$$

In[3]:= **r = Sqrt[(s1 - s0) p + (s2 - s0) q + s0];**

In[4]:= **dW = (1 / 2) * (Br^2 / mu r + Bz^2 / mu z)**

Out[4]=
$$\frac{1}{2} \left(\frac{(b_1 (-\phi_0 + \phi_1) + b_2 (-\phi_0 + \phi_2))^2}{d^2 \pi^2 \mu z} + \frac{(c_1 (-\phi_0 + \phi_1) + c_2 (-\phi_0 + \phi_2))^2}{4 d^2 \pi^2 (s_0 + p (-s_0 + s_1) + q (-s_0 + s_2)) \mu r} \right)$$

We want to integrate $2 \pi r dW dr dz$. Note that $dr = ds \cdot (dr/ds) = ds \cdot D[\text{Sqrt}[s], s] = ds / (2 \cdot \text{Sqrt}[s]) = ds / (2 \cdot r)$ so that $2 \pi r dW dr dz = \pi dW ds dz$. To integrate over p and q , turns into $\pi dW d p dq$

In[5]:= **X = FullSimplify[Integrate[FullSimplify[Pi dW d], {p, 0, 1 - q}]] [[1]]**

Out[5]=
$$\frac{-\frac{4 (-1+q) (b_1 (\phi_0-\phi_1)+b_2 (\phi_0-\phi_2))^2}{\mu z} + \frac{(c_1 (\phi_0-\phi_1)+c_2 (\phi_0-\phi_2))^2 (\text{Log}[s_0-q s_0+q s_2]-\text{Log}[s_1-q s_1+q s_2])}{(s_0-s_1) \mu r}}{8 d \pi}$$

In[6]:= **Y = Integrate[X, {q, 0, 1}] [[1]]**

$$\text{Out[6]} = \frac{\frac{2 (b_1 (\phi_0 - \phi_1) + b_2 (\phi_0 - \phi_2))^2}{\mu z} + \frac{(c_1 (\phi_0 - \phi_1) + c_2 (\phi_0 - \phi_2))^2 (s_0 (-s_1 + s_2) \text{Log}[s_0] + s_1 (s_0 - s_2) \text{Log}[s_1] + (-s_0 + s_1) s_2 \text{Log}[s_2])}{(s_0 - s_1) (s_0 - s_2) (-s_1 + s_2) \mu r}}{8 d \pi}$$

Y is an expression for the stored energy in an element. Differentiate with respect to the nodal values of flux to get the element matrices

In[7]:= **m = {D[Y, ϕ0], D[Y, ϕ1], D[Y, ϕ2]}**

$$\text{Out[7]} = \left\{ \frac{1}{8 d \pi} \left(\frac{4 (b_1 + b_2) (b_1 (\phi_0 - \phi_1) + b_2 (\phi_0 - \phi_2))}{\mu z} + \frac{(2 (c_1 + c_2) (c_1 (\phi_0 - \phi_1) + c_2 (\phi_0 - \phi_2)) (s_0 (-s_1 + s_2) \text{Log}[s_0] + s_1 (s_0 - s_2) \text{Log}[s_1] + (-s_0 + s_1) s_2 \text{Log}[s_2]))}{((s_0 - s_1) (s_0 - s_2) (-s_1 + s_2) \mu r)} \right), \right. \\ \left. - \frac{4 b_1 (b_1 (\phi_0 - \phi_1) + b_2 (\phi_0 - \phi_2))}{\mu z} - \frac{2 c_1 (c_1 (\phi_0 - \phi_1) + c_2 (\phi_0 - \phi_2)) (s_0 (-s_1 + s_2) \text{Log}[s_0] + s_1 (s_0 - s_2) \text{Log}[s_1] + (-s_0 + s_1) s_2 \text{Log}[s_2])}{(s_0 - s_1) (s_0 - s_2) (-s_1 + s_2) \mu r}, \right. \\ \left. - \frac{4 b_2 (b_1 (\phi_0 - \phi_1) + b_2 (\phi_0 - \phi_2))}{\mu z} - \frac{2 c_2 (c_1 (\phi_0 - \phi_1) + c_2 (\phi_0 - \phi_2)) (s_0 (-s_1 + s_2) \text{Log}[s_0] + s_1 (s_0 - s_2) \text{Log}[s_1] + (-s_0 + s_1) s_2 \text{Log}[s_2])}{(s_0 - s_1) (s_0 - s_2) (-s_1 + s_2) \mu r} \right\}$$

In[27]:= $\mathbf{M} = \{D[m, \phi_0], D[m, \phi_1], D[m, \phi_2]\}$

$$\text{Out[27]} = \left\{ \left\{ \frac{\frac{4(b_1+b_2)^2}{\mu z} + \frac{2(c_1+c_2)^2(s_0(-s_1+s_2)\text{Log}[s_0]+s_1(s_0-s_2)\text{Log}[s_1]+(-s_0+s_1)s_2\text{Log}[s_2])}{(s_0-s_1)(s_0-s_2)(-s_1+s_2)\mu r}}{8d\pi}, \right. \right. \\ \left. \frac{-\frac{4b_1(b_1+b_2)}{\mu z} - \frac{2c_1(c_1+c_2)(s_0(-s_1+s_2)\text{Log}[s_0]+s_1(s_0-s_2)\text{Log}[s_1]+(-s_0+s_1)s_2\text{Log}[s_2])}{(s_0-s_1)(s_0-s_2)(-s_1+s_2)\mu r}}{8d\pi}, \right. \\ \left. \frac{-\frac{4b_2(b_1+b_2)}{\mu z} - \frac{2c_2(c_1+c_2)(s_0(-s_1+s_2)\text{Log}[s_0]+s_1(s_0-s_2)\text{Log}[s_1]+(-s_0+s_1)s_2\text{Log}[s_2])}{(s_0-s_1)(s_0-s_2)(-s_1+s_2)\mu r}}{8d\pi} \right\}, \\ \left\{ \frac{-\frac{4b_1(b_1+b_2)}{\mu z} - \frac{2c_1(c_1+c_2)(s_0(-s_1+s_2)\text{Log}[s_0]+s_1(s_0-s_2)\text{Log}[s_1]+(-s_0+s_1)s_2\text{Log}[s_2])}{(s_0-s_1)(s_0-s_2)(-s_1+s_2)\mu r}}{8d\pi}, \right. \\ \frac{\frac{4b_1^2}{\mu z} + \frac{2c_1^2(s_0(-s_1+s_2)\text{Log}[s_0]+s_1(s_0-s_2)\text{Log}[s_1]+(-s_0+s_1)s_2\text{Log}[s_2])}{(s_0-s_1)(s_0-s_2)(-s_1+s_2)\mu r}}{8d\pi}, \\ \left. \frac{\frac{4b_1b_2}{\mu z} + \frac{2c_1c_2(s_0(-s_1+s_2)\text{Log}[s_0]+s_1(s_0-s_2)\text{Log}[s_1]+(-s_0+s_1)s_2\text{Log}[s_2])}{(s_0-s_1)(s_0-s_2)(-s_1+s_2)\mu r}}{8d\pi} \right\}, \\ \left\{ \frac{-\frac{4b_2(b_1+b_2)}{\mu z} - \frac{2c_2(c_1+c_2)(s_0(-s_1+s_2)\text{Log}[s_0]+s_1(s_0-s_2)\text{Log}[s_1]+(-s_0+s_1)s_2\text{Log}[s_2])}{(s_0-s_1)(s_0-s_2)(-s_1+s_2)\mu r}}{8d\pi}, \right. \\ \frac{\frac{4b_1b_2}{\mu z} + \frac{2c_1c_2(s_0(-s_1+s_2)\text{Log}[s_0]+s_1(s_0-s_2)\text{Log}[s_1]+(-s_0+s_1)s_2\text{Log}[s_2])}{(s_0-s_1)(s_0-s_2)(-s_1+s_2)\mu r}}{8d\pi}, \\ \left. \frac{\frac{4b_2^2}{\mu z} + \frac{2c_2^2(s_0(-s_1+s_2)\text{Log}[s_0]+s_1(s_0-s_2)\text{Log}[s_1]+(-s_0+s_1)s_2\text{Log}[s_2])}{(s_0-s_1)(s_0-s_2)(-s_1+s_2)\mu r}}{8d\pi} \right\} \right\}$$

Split into element matrices for the r - and z - directions

In[9]:= $\mathbf{Mz} = \text{FullSimplify}[\text{Limit}[\mathbf{M}, \mu r \rightarrow \text{Infinity}]] /. b_1 + b_2 \rightarrow -b_0;$
 $\text{MatrixForm}[\mathbf{Mz}]$

Out[10]//MatrixForm=

$$\begin{pmatrix} \frac{b_0^2}{2d\pi\mu z} & \frac{b_0b_1}{2d\pi\mu z} & \frac{b_0b_2}{2d\pi\mu z} \\ \frac{b_0b_1}{2d\pi\mu z} & \frac{b_1^2}{2d\pi\mu z} & \frac{b_1b_2}{2d\pi\mu z} \\ \frac{b_0b_2}{2d\pi\mu z} & \frac{b_1b_2}{2d\pi\mu z} & \frac{b_2^2}{2d\pi\mu z} \end{pmatrix}$$

In[24]:= **Mr = FullSimplify[Limit[M, $\mu z \rightarrow \text{Infinity}$]] /. c1 + c2 \rightarrow -c0**

$$\text{Out[24]= } \left\{ \left\{ \frac{c_0^2 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{4 d \pi (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r}, \right. \right. \\ \left. \frac{c_0 c_1 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{4 d \pi (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r}, \right. \\ \left. \frac{c_0 c_2 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{4 d \pi (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r} \right\}, \\ \left\{ \frac{c_0 c_1 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{4 d \pi (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r}, \right. \\ \left. \frac{c_1^2 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{4 d \pi (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r}, \right. \\ \left. \frac{c_1 c_2 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{4 d \pi (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r} \right\}, \\ \left\{ \frac{c_0 c_2 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{4 d \pi (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r}, \right. \\ \left. \frac{c_1 c_2 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{4 d \pi (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r}, \right. \\ \left. \frac{c_2^2 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{4 d \pi (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r} \right\} \right\}$$

However, I want to implement the code in terms of vector potential, A, for various reasons. We can convert the element matrices for use with nodal values of A with the matrix T, i.e. $\{\phi_0, \phi_1, \phi_2\} = T \cdot \{A_0, A_1, A_3\}$ where T is:

In[13]:= **T = 2 Pi DiagonalMatrix[{r0, r1, r2}]**

Out[13]= $\{\{2 \pi r_0, 0, 0\}, \{0, 2 \pi r_1, 0\}, \{0, 0, 2 \pi r_2\}\}$

Z - direction element matrix for use with vector potential is :

In[14]:= **Mza = T.Mz.T;**
MatrixForm[Mza]

Out[15]//MatrixForm=

$$\begin{pmatrix} \frac{2 b_0^2 \pi r_0^2}{d \mu z} & \frac{2 b_0 b_1 \pi r_0 r_1}{d \mu z} & \frac{2 b_0 b_2 \pi r_0 r_2}{d \mu z} \\ \frac{2 b_0 b_1 \pi r_0 r_1}{d \mu z} & \frac{2 b_1^2 \pi r_1^2}{d \mu z} & \frac{2 b_1 b_2 \pi r_1 r_2}{d \mu z} \\ \frac{2 b_0 b_2 \pi r_0 r_2}{d \mu z} & \frac{2 b_1 b_2 \pi r_1 r_2}{d \mu z} & \frac{2 b_2^2 \pi r_2^2}{d \mu z} \end{pmatrix}$$

R - direction element matrix for use with vector potential is :

In[25]= **Mra = T.Mr.T**

$$\text{Out[25]= } \left\{ \left\{ \frac{c_0^2 \pi r_0^2 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{d (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r}, \right. \right. \\ \left. \frac{c_0 c_1 \pi r_0 r_1 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{d (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r}, \right. \\ \left. \frac{c_0 c_2 \pi r_0 r_2 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{d (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r} \right\}, \\ \left\{ \frac{c_0 c_1 \pi r_0 r_1 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{d (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r}, \right. \\ \left. \frac{c_1^2 \pi r_1^2 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{d (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r}, \right. \\ \left. \frac{c_1 c_2 \pi r_1 r_2 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{d (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r} \right\}, \\ \left\{ \frac{c_0 c_2 \pi r_0 r_2 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{d (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r}, \right. \\ \left. \frac{c_1 c_2 \pi r_1 r_2 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{d (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r}, \right. \\ \left. \frac{c_2^2 \pi r_2^2 (s_0 (s_1 - s_2) \text{Log}[s_0] + s_1 (-s_0 + s_2) \text{Log}[s_1] + (s_0 - s_1) s_2 \text{Log}[s_2])}{d (s_0 - s_1) (s_0 - s_2) (s_1 - s_2) \mu r} \right\} \right\}$$

For the r-direction case, we can do a bit of extra work to define matrices when one of the nodes lies on $r=0$

In[21]= **MatrixForm[Limit[Mra, $s_0 \rightarrow 0$] /. $r_0 \rightarrow 0$]**

$$\text{Out[21]//MatrixForm= } \begin{pmatrix} 0 & 0 & 0 \\ 0 & \frac{c_1^2 \pi r_1^2 (\text{Log}[s_1] - \text{Log}[s_2])}{d (s_1 - s_2) \mu r} & \frac{c_1 c_2 \pi r_1 r_2 (\text{Log}[s_1] - \text{Log}[s_2])}{d (s_1 - s_2) \mu r} \\ 0 & \frac{c_1 c_2 \pi r_1 r_2 (\text{Log}[s_1] - \text{Log}[s_2])}{d (s_1 - s_2) \mu r} & \frac{c_2^2 \pi r_2^2 (\text{Log}[s_1] - \text{Log}[s_2])}{d (s_1 - s_2) \mu r} \end{pmatrix}$$

In[22]= **MatrixForm[Limit[Mra, $s_1 \rightarrow 0$] /. $r_1 \rightarrow 0$]**

$$\text{Out[22]//MatrixForm= } \begin{pmatrix} \frac{c_0^2 \pi r_0^2 (\text{Log}[s_0] - \text{Log}[s_2])}{d (s_0 - s_2) \mu r} & 0 & \frac{c_0 c_2 \pi r_0 r_2 (\text{Log}[s_0] - \text{Log}[s_2])}{d (s_0 - s_2) \mu r} \\ 0 & 0 & 0 \\ \frac{c_0 c_2 \pi r_0 r_2 (\text{Log}[s_0] - \text{Log}[s_2])}{d (s_0 - s_2) \mu r} & 0 & \frac{c_2^2 \pi r_2^2 (\text{Log}[s_0] - \text{Log}[s_2])}{d (s_0 - s_2) \mu r} \end{pmatrix}$$

In[23]= **MatrixForm[Limit[Mra, $s_2 \rightarrow 0$] /. $r_2 \rightarrow 0$]**

$$\text{Out[23]//MatrixForm= } \begin{pmatrix} \frac{c_0^2 \pi r_0^2 (\text{Log}[s_0] - \text{Log}[s_1])}{d (s_0 - s_1) \mu r} & \frac{c_0 c_1 \pi r_0 r_1 (\text{Log}[s_0] - \text{Log}[s_1])}{d (s_0 - s_1) \mu r} & 0 \\ \frac{c_0 c_1 \pi r_0 r_1 (\text{Log}[s_0] - \text{Log}[s_1])}{d (s_0 - s_1) \mu r} & \frac{c_1^2 \pi r_1^2 (\text{Log}[s_0] - \text{Log}[s_1])}{d (s_0 - s_1) \mu r} & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

For the two - nodes - on - zero case, the Mra matrix is all zeros--there's no stored energy in the R direc-

tion because $Br=0$